

RELATIVELY INDEPENDENT JOININGS AND SUBSYSTEMS OF W^* -DYNAMICAL SYSTEMS

ROCCO DUVENHAGE

ABSTRACT. Relatively independent joinings of W^* -dynamical systems are constructed. This is intimately related to subsystems of W^* -dynamical systems, and therefore we also study general properties of subsystems, in particular fixed point subsystems and compact subsystems. This allows us to obtain characterizations of weak mixing and relative ergodicity, as well as of certain compact subsystems, in terms of joinings.

1. INTRODUCTION

In the study of the general mathematical structure of quantum dynamical systems and quantum statistical mechanics, the operator algebraic approach has proven very valuable. In particular the framework of von Neumann algebras, in which case the dynamical system is called a W^* -dynamical system (see Section 2 for a precise definition), provides a natural arena for ergodic theory which extends the classical measure theoretic framework. Refer to [6, 7] for an account of many aspects of these topics.

The first aim of this paper is to study the extension of the concept of a relatively independent joining of two dynamical systems in classical ergodic theory to the noncommutative framework of W^* -dynamical systems. It is essentially a generalized way of forming the tensor product of two systems which takes into account a common subsystem of the systems. A clear exposition of the classical case can be found in [15, Chapter 6].

Since the early works [14, 30], joinings have been a useful tool in classical ergodic theory (again see [15]). This paper, building on [9, 10], is part of a programme to systematically develop the theory of joinings in the noncommutative case. The goal is to eventually have a similarly useful tool for noncommutative dynamical systems.

Joinings have in fact already gradually found some use in noncommutative dynamical systems, in particular related to dynamical entropy [31] where a special case of joinings appears (also see [25, Section 5.1]).

Date: 2010-5-15.

2000 Mathematics Subject Classification. Primary 46L55.

Key words and phrases. W^* -dynamical systems; relatively independent joinings; subsystems.

Early work related to joinings in a noncommutative setting goes back to at least [4], where disjointness was studied in the context of a noncommutative extension of topological dynamics (see in particular [4, Definition 2.9 and Section 5]).

Since subsystems form an integral part of relatively independent joinings, we also study the properties of particular subsystems, namely the fixed point subsystem and the compact subsystem generated by the eigenoperators of the dynamics. These subsystems are related to the properties of W^* -dynamical systems, namely to ergodicity and weak mixing respectively. For subsystems more generally we will see in Section 2 and onward that the modular group of the state of the W^* -dynamical system in question plays an essential role in the definition and application of subsystems relevant to relatively independent joinings. This fits in very naturally with the physically relevant case in which the dynamics itself is given by the modular group of the state, but we will see that the framework is much more widely applicable. Note that subsystems have already proven useful in the study of noncommutative dynamical systems; see for example [3, Sections 3 and 4] and [28, Definition 2.9 and Lemma 2.10] for recent work.

After summarizing the basic framework in Section 2, we construct relatively independent joinings of two W^* -dynamical systems in Section 3 and derive a few useful facts regarding them. In Sections 4 and 5 we then study relative ergodicity and compact subsystems respectively, to illustrate how relatively independent joinings fit into the theory of W^* -dynamical systems. In the latter two sections the fixed point subsystem and compact subsystem generated by the eigenoperators respectively play a central role. In Section 6 we conclude the paper by studying further properties of this compact subsystem when the W^* -dynamical system is ergodic.

2. BASIC DEFINITIONS, NOTATIONS AND BACKGROUND

We use the same basic definitions as in [9, 10]. For convenience we summarize them here, along with some additional definitions. Simultaneously this fixes notations that will be used throughout the rest of the paper. Some related background material, in particular regarding Tomita-Takesaki theory (or modular theory) is also discussed. A general notation that we use often is $B(X)$ to denote the space of all bounded linear operators $X \rightarrow X$ on a normed space X . The identity element of a group will be indicated by 1. In the remainder of this paper W^* -dynamical systems are referred to simply as “systems” and they are defined as follows:

Definition 2.1. A *system* $\mathbf{A} = (A, \mu, \alpha)$ consists of a faithful normal state μ on a (necessarily σ -finite) von Neumann algebra A , and a representation $\alpha : G \rightarrow \text{Aut}(A) : g \mapsto \alpha_g$ of an arbitrary group G as

$*$ -automorphisms of A , such that $\mu \circ \alpha_g = \mu$ for all g . We call the system \mathbf{A} an *identity system* if $\alpha_g = \text{id}_A$ for all g where $\text{id}_A : A \rightarrow A$ is the identity mapping, while we call it *trivial* if $A = \mathbb{C}1_A$ where 1_A (often denoted simply as 1) is the unit of A .

In the rest of the paper the symbols \mathbf{A} , \mathbf{B} , \mathbf{F} and \mathbf{R} will denote systems (A, μ, α) , (B, ν, β) , (F, λ, φ) and (R, ψ, ρ) respectively, all making use of actions of the same group G .

Definition 2.2. A *joining* of \mathbf{A} and \mathbf{B} is a state ω (i.e. a positive linear functional with $\omega(1) = 1$) on the algebraic tensor product $A \odot B$ such that $\omega(a \otimes 1_B) = \mu(a)$, $\omega(1_A \otimes b) = \nu(b)$ and $\omega \circ (\alpha_g \odot \beta_g) = \omega$ for all $a \in A$, $b \in B$ and $g \in G$. The set of all joinings of \mathbf{A} and \mathbf{B} is denoted by $J(\mathbf{A}, \mathbf{B})$. We call \mathbf{A} *disjoint* from \mathbf{B} when $J(\mathbf{A}, \mathbf{B}) = \{\mu \odot \nu\}$.

We will also have occasion to use a more general concept, namely if A and B are von Neumann algebras with faithful normal states μ and ν respectively, then a *coupling* of (A, μ) and (B, ν) is a state ω on $A \odot B$ such that $\omega(a \otimes 1_B) = \mu(a)$ and $\omega(1_A \otimes b) = \nu(b)$.

The modular group of a faithful normal state μ on a von Neumann algebra A will be denoted by σ^μ , and its elements by σ_t^μ for every $t \in \mathbb{R}$. The cyclic representation of A obtained from μ by the GNS construction will be denoted by $(H_\mu, \pi_\mu, \Omega_\mu)$; this notation will in fact also be used in the case of an arbitrary state on a unital $*$ -algebra. The associated modular conjugation will be denoted by J_μ , and we let

$$j_\mu : B(H_\mu) \rightarrow B(H_\mu) : a \mapsto J_\mu a^* J_\mu$$

for which we note that $j_\mu^{-1} = j_\mu$. We will also use the notation

$$\gamma_\mu : A \rightarrow H_\mu : a \mapsto \pi_\mu(a)\Omega_\mu$$

even in the case of unital $*$ -algebras.

The dynamics α of a system \mathbf{A} can be represented by a unitary group U on H_μ defined by extending

$$U_g \gamma_\mu(a) := \gamma_\mu(\alpha_g(a)).$$

It satisfies

$$U_g \pi_\mu(a) U_g^* = \pi_\mu(\alpha_g(a))$$

for all $g \in G$; also see [6, Corollary 2.3.17]. The unitary representation of β will be denoted by V .

Definition 2.3. We call \mathbf{F} a *subsystem* of \mathbf{A} if there exists an injective unital $*$ -homomorphism ζ of F onto a von Neumann subalgebra of A such that $\mu \circ \zeta = \lambda$ and $\alpha_g \circ \zeta = \zeta \circ \varphi_g$ for all $g \in G$. If $\zeta(F)$ is invariant under σ^μ , i.e. $\sigma_t^\mu(\zeta(F)) = \zeta(F)$ for all $t \in \mathbb{R}$, then \mathbf{F} is called a *modular subsystem* of \mathbf{A} . If furthermore $\zeta : F \rightarrow A$ is surjective, then we say that ζ is an *isomorphism* of dynamical systems, and the systems \mathbf{A} and \mathbf{F} are *isomorphic*.

The symbol ζ will be used uniformly for this purpose. We call it the *imbedding* of \mathbf{F} into \mathbf{A} . Without loss one could assume ζ to be an inclusion, so $F \subset A$, and we will often do so. When the system \mathbf{F} is a subsystem of \mathbf{B} , the notation $\eta : F \rightarrow B$ instead of ζ will be used. If \mathbf{F} is a (modular) subsystem of both \mathbf{A} and \mathbf{B} , then we call it a *common* (modular) subsystem of \mathbf{A} and \mathbf{B} .

Important nontrivial examples of modular subsystems will be studied in Sections 4 and 5. In the meantime we note that in classical ergodic theory all subsystems (also known as “factors”) are of course modular. More generally, if the state μ of our system \mathbf{A} happens to be a trace (i.e. $\mu(ab) = \mu(ba)$ for all $a, b \in A$), then again all subsystems are modular. Also, in the physically relevant situation where the dynamics α of \mathbf{A} is the modular group of μ , then any subsystem of \mathbf{A} is automatically modular.

A standard fact from Tomita-Takesaki theory related to modular subsystems is the following (a proof of which is contained in [35, Section 10.2] for example):

Lemma 2.4. *Let μ be a faithful normal state on a von Neumann algebra A . Let F be a von Neumann subalgebra of A , invariant under σ^μ . Setting $H_F := \overline{\gamma_\mu(F)}$, we have $J_\mu H_F = H_F$.*

When we construct relatively independent joinings of \mathbf{A} and \mathbf{B} over a common modular subsystem \mathbf{F} in the next section, one of the key tricks (also used in [12]) is to work with the “commutant” $\tilde{\mathbf{B}}$ of \mathbf{B} . Given a system \mathbf{B} , this means the following: We set $\tilde{B} := \pi_\nu(B)'$ and then carry the state and dynamics of \mathbf{B} over to \tilde{B} in a natural way using j_ν by defining a state $\tilde{\nu}$ and $*$ -automorphism $\tilde{\beta}_g$ on \tilde{B} by $\tilde{\nu}(b) := \nu \circ \pi_\nu^{-1} \circ j_\nu(b)$ and $\tilde{\beta}_g(b) := j_\nu \circ \pi_\nu \circ \beta_g \circ \pi_\nu^{-1} \circ j_\nu(b)$ for all $g \in G$. From Tomita-Takesaki theory one has that $V_g J_\nu = J_\nu V_g$ (see [9, Construction 3.4]). It follows that

$$\tilde{\nu}(b) = \langle \Omega_\nu, b \Omega_\nu \rangle$$

and

$$\tilde{\beta}_g(b) = V_g b V_g^*$$

for all $b \in \tilde{B}$ and $g \in G$. In particular the latter tells us that the unitary representation of $\tilde{\beta}$ is the same as that of β , namely V .

3. RELATIVELY INDEPENDENT JOININGS

Throughout this section we consider two systems \mathbf{A} and \mathbf{B} which have a common modular subsystem \mathbf{F} . We are going to construct the relatively independent joining of \mathbf{A} and \mathbf{B} over \mathbf{F} . More precisely it will be a joining of \mathbf{A} and $\tilde{\mathbf{B}}$. Without loss we can assume ζ to be an inclusion, so $F \subset A$.

Since \mathbf{F} is a modular subsystem of \mathbf{B} , we obtain a modular subsystem $\tilde{\mathbf{F}} = (\tilde{F}, \tilde{\lambda}, \tilde{\varphi})$ of $\tilde{\mathbf{B}}$ in the following way (as is easily checked): Set

$$\tilde{F} := j_\nu \circ \pi_\nu \circ \eta(F) \subset \tilde{B}$$

and let $\tilde{\lambda} := \tilde{\nu}|_{\tilde{F}}$ and $\tilde{\varphi}_g := \tilde{\beta}_g|_{\tilde{F}}$ for all $g \in G$.

Since F and \tilde{F} are invariant under the modular groups of μ and $\tilde{\nu}$ respectively we know by Tomita-Takesaki theory (see for example [36, Theorem IX.4.2]) that we have unique conditional expectations

$$D : A \rightarrow F$$

and

$$\tilde{D} : \tilde{B} \rightarrow \tilde{F}$$

such that $\lambda \circ D = \mu$ and $\tilde{\lambda} \circ \tilde{D} = \tilde{\nu}$. (This corresponds to disintegrations in the classical case; see [15, Section 6.1].)

Setting $H_\eta := \overline{\gamma_\nu(\eta(F))}$, we obtain a well-defined unital $*$ -homomorphism

$$\pi_\eta : F \rightarrow B(H_\eta) : a \mapsto \pi_\nu(\eta(a))|_{H_\eta}$$

and one can check that $(H_\eta, \pi_\eta, \Omega_\nu)$ is a cyclic representation of (F, λ) and hence unitarily equivalent [6, Theorem 2.3.16] to $(H_\lambda, \pi_\lambda, \Omega_\lambda)$. In other words there is a unique unitary operator $u_\eta : H_\lambda \rightarrow H_\eta$ such that $u_\eta \Omega_\lambda = \Omega_\nu$ and $u_\eta^* \pi_\eta(a) u_\eta = \pi_\lambda(a)$ for all $a \in F$. From Lemma 2.4, for any $a \in \tilde{F}$ we have $a H_\eta \subset H_\eta$, and therefore we obtain a well-defined injective $*$ -homomorphism

$$\kappa_\eta : \tilde{F} \rightarrow B(H_\lambda) : a \mapsto u_\eta^*(a|_{H_\eta}) u_\eta.$$

It is then easily shown from the definition of \tilde{F} that $\kappa_\eta(\tilde{F}) \subset \pi_\lambda(F)'$, and hence

$$\delta : \pi_\lambda(F) \odot \kappa_\eta(\tilde{F}) \rightarrow B(H_\lambda)$$

defined as the linear extension of $\pi_\lambda(F) \times \kappa_\eta(\tilde{F}) \rightarrow B(H_\lambda) : (a, b) \mapsto ab$ is a well-defined unital $*$ -homomorphism. We now introduce the *diagonal* state of λ as the state

$$\Delta_\lambda : F \odot \tilde{F} \rightarrow \mathbb{C}$$

defined by

$$\Delta_\lambda(c) := \langle \Omega_\lambda, \delta \circ (\pi_\lambda \odot \kappa_\eta)(c) \Omega_\lambda \rangle$$

for all $c \in F \odot \tilde{F}$. This enables us to define a linear functional $\mu \odot_\lambda \tilde{\nu}$ on $A \odot \tilde{B}$ by

$$\mu \odot_\lambda \tilde{\nu} := \Delta_\lambda \circ (D \odot \tilde{D}).$$

(If we did not assume ζ to be an inclusion mapping, one would simply replace D by $\zeta^{-1} \circ D$.)

This completes the basic construction. The next step is to show that $\mu \odot_\lambda \tilde{\nu}$ is a joining of \mathbf{A} and $\tilde{\mathbf{B}}$. In fact we will show a bit more in the next proposition. We first define a special class of joinings of \mathbf{A} and $\tilde{\mathbf{B}}$:

Definition 3.1. Any $\omega \in J(\mathbf{A}, \tilde{\mathbf{B}})$ with $\omega|_{F \odot \tilde{F}} = \Delta_\lambda$ is called a *joining of \mathbf{A} and $\tilde{\mathbf{B}}$ over \mathbf{F}* . The set of all such ω is denoted $J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$.

Proposition 3.2. If \mathbf{F} is a common modular subsystem of \mathbf{A} and \mathbf{B} , then $\mu \odot_\lambda \tilde{\nu} \in J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$.

Proof. From the uniqueness of D it follows that $\varphi_g^{-1} \circ D \circ \alpha_g = D$ and therefore $D \circ \alpha_g = \varphi_g \circ D$. Similarly $\tilde{D} \circ \tilde{\beta}_g = \tilde{\varphi}_g \circ \tilde{D}$. Denoting the unitary representation of φ on H_λ by W , note that for $a \in F$ we have $u_\eta^* V_g|_{H_\eta} u_\eta \pi_\lambda(a) u_\eta^* V_g|_{H_\eta} u_\eta = u_\eta^* \pi_\nu(\beta_g(\eta(a))) u_\eta = \pi_\lambda(\varphi_g(a)) = W_g \pi_\lambda(a) W_g^*$. Letting this act on the cyclic vector Ω_λ , we obtain $u_\eta^* V_g|_{H_\eta} u_\eta = W_g$. Therefore $\kappa_\eta(\tilde{\varphi}_g(b)) = W_g \kappa_\eta(b) W_g^*$ for $b \in \tilde{F}$.

For $a \in A$ and $b \in \tilde{B}$ it follows that

$$\begin{aligned} & \mu \odot_\lambda \tilde{\nu}(\alpha_g \odot \tilde{\beta}_g(a \otimes b)) \\ &= \langle \Omega_\lambda, \pi_\lambda(\varphi_g(D(a))) \kappa_\eta(\tilde{\varphi}_g(\tilde{D}(b))) \Omega_\lambda \rangle \\ &= \langle \Omega_\lambda, W_g \pi_\lambda(D(a)) W_g^* W_g \kappa_\eta(\tilde{D}(b)) W_g^* \Omega_\lambda \rangle \\ &= \mu \odot_\lambda \tilde{\nu}(a \otimes b) \end{aligned}$$

and therefore $\mu \odot_\lambda \tilde{\nu} \circ (\alpha_g \odot \tilde{\beta}_g) = \mu \odot_\lambda \tilde{\nu}$ as required. Since D and \tilde{D} are conditional expectations, it follows that $D \odot \tilde{D}$ is positive (see [35, p. 119] for example). It is now easily seen that $\mu \odot_\lambda \tilde{\nu} \in J(\mathbf{A}, \tilde{\mathbf{B}})$. From the definition of $\mu \odot_\lambda \tilde{\nu}$ it then immediately follows that $\mu \odot_\lambda \tilde{\nu} \in J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$. \square

Definition 3.3. The joining $\mu \odot_\lambda \tilde{\nu}$ of \mathbf{A} and $\tilde{\mathbf{B}}$ is called the *relatively independent joining of \mathbf{A} and $\tilde{\mathbf{B}}$ over \mathbf{F}* . If $J_\lambda(\mathbf{A}, \tilde{\mathbf{B}}) = \{\mu \odot_\lambda \tilde{\nu}\}$, then \mathbf{A} and $\tilde{\mathbf{B}}$ are called *disjoint over \mathbf{F}* .

In the remainder of this section we study simple but useful properties of this relatively independent joining. These properties will be used in subsequent sections of the paper.

Proposition 3.4. If $F = \mathbb{C}$, it follows that $\mu \odot_\lambda \tilde{\nu} = \mu \odot \tilde{\nu}$ and $J_\lambda(\mathbf{A}, \tilde{\mathbf{B}}) = J(\mathbf{A}, \tilde{\mathbf{B}})$.

Proof. This follows directly from the definitions above. \square

Proposition 3.5. Let \mathbf{F} and \mathbf{R} both be common modular subsystems of \mathbf{A} and \mathbf{B} , and \mathbf{F} a subsystem of \mathbf{R} , with inclusions $F \subset R \subset A$ giving the corresponding imbeddings of the subsystems into \mathbf{A} , while $\theta : R \rightarrow B$ and $\eta = \theta|_F$ give the imbeddings into \mathbf{B} . Then $J_\psi(\mathbf{A}, \tilde{\mathbf{B}}) \subset J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$. Furthermore, if $\mu \odot_\psi \tilde{\nu}|_{R \odot \tilde{R}} = \mu \odot_\lambda \tilde{\nu}|_{R \odot \tilde{R}}$, then $R = F$.

Proof. For the first part we only need to show that $\Delta_\psi|_{F \odot \tilde{F}} = \Delta_\lambda$. Given $(H_\psi, \pi_\psi, \Omega_\psi)$, we can assume that $(H_\lambda, \pi_\lambda, \Omega_\lambda)$ is given by $H_\lambda = \overline{\gamma_\psi(\tilde{F})}$, $\pi_\lambda(a) = \pi_\psi(a)|_{H_\lambda}$ for $a \in F$, and $\Omega_\lambda = \Omega_\psi$, without changing Δ_λ , because of unitary equivalence between different cyclic representations of the same state. We can also define $\kappa_\theta : \tilde{R} \rightarrow B(H_\psi)$ in the same way as κ_η above, with the corresponding unitary operator denoted by u_θ instead of u_η . Note that $u_\theta \gamma_\lambda(a) = \pi_\theta(a) \Omega_\nu = \pi_\eta(a) \Omega_\nu = u_\eta \gamma_\lambda(a)$ for $a \in F$, from which follows that $\kappa_\theta(b)|_{H_\lambda} = \kappa_\eta(b)$ for $b \in \tilde{F}$, hence $\Delta_\psi|_{F \odot \tilde{F}} = \Delta_\lambda$.

Assuming $\mu \odot_\psi \tilde{\nu}|_{R \odot \tilde{R}} = \mu \odot_\lambda \tilde{\nu}|_{R \odot \tilde{R}}$, we have (in H_ψ)

$$\begin{aligned} \langle \pi_\psi(a) \Omega_\psi, \kappa_\theta(b) \Omega_\psi \rangle &= \mu \odot_\psi \tilde{\nu}(a^* \otimes b) = \mu \odot_\lambda \tilde{\nu}(a^* \otimes b) \\ &= \mu \odot_\lambda \tilde{\nu}(a^* \otimes \tilde{D}(b)) = \mu \odot_\psi \tilde{\nu}(a^* \otimes \tilde{D}(b)) \\ &= \langle \pi_\psi(a) \Omega_\psi, \kappa_\theta(\tilde{D}(b)) \Omega_\psi \rangle \end{aligned}$$

for any $a \in R$ and $b \in \tilde{R}$ (with \tilde{R} of course defined in an analogous way to \tilde{F}). Since Ω_ψ is cyclic for π_ψ , and therefore separating for $\kappa_\theta(\tilde{R}) \subset \pi_\psi(R)'$, it follows that $\kappa_\theta(b) = \kappa_\theta(\tilde{D}(b))$. Therefore $b|_{H_\theta} = \tilde{D}(b)|_{H_\theta}$ where $H_\theta := \overline{\gamma_\nu(\theta(R))}$, and hence $b \Omega_\nu = \tilde{D}(b) \Omega_\nu$. Since Ω_ν is separating for \tilde{B} we conclude that $b = \tilde{D}(b) \in \tilde{F}$ and therefore $\tilde{R} = \tilde{F}$. It follows that $R = F$ as required. \square

Finally we consider a Hilbert space characterization of relatively independent joinings. We will work in the following setting: Let ω be a coupling of (A, μ) and $(\tilde{B}, \tilde{\nu})$ as defined in Section 2. Then we can “imbed” $(H_\mu, \pi_\mu, \Omega_\mu)$ and $(H_{\tilde{\nu}}, \pi_{\tilde{\nu}}, \Omega_{\tilde{\nu}})$ into $(H_\omega, \pi_\omega, \Omega_\omega)$ in a natural way (see [9, Construction 2.3] for explicit details; also note that $(H_{\tilde{\nu}}, \pi_{\tilde{\nu}}, \Omega_{\tilde{\nu}})$ is unitarily equivalent to $(H_\nu, \iota_{\tilde{B}}, \Omega_\nu)$ but not necessarily equal). In particular H_μ and $H_{\tilde{\nu}}$ are then subspaces of H_ω and the corresponding cyclic vectors are equal to Ω_ω . We set $H'_\lambda = \overline{\gamma_\mu(F)} \subset H_\mu$ and $H_{\tilde{\lambda}} = \overline{\gamma_{\tilde{\nu}}(\tilde{F})} \subset H_{\tilde{\nu}}$,

We are particularly interested in the case where $\omega|_{F \odot \tilde{F}} = \Delta_\lambda$. If this condition is satisfied one has $H'_\lambda = H_{\tilde{\lambda}}$ as follows: H'_λ and $H_{\tilde{\lambda}}$ are both contained $H_\odot := \overline{\gamma_\omega(F \odot \tilde{F})}$. Consider $x \in H_\odot \ominus H'_\lambda$ and a sequence $b_n \in F \odot \tilde{F}$ such that $\gamma_\omega(b_n) \rightarrow x$, then for every $a \in F$ it follows from $\omega|_{F \odot \tilde{F}} = \Delta_\lambda$ that $0 = \langle \gamma_\mu(a), x \rangle = \langle \gamma_\lambda(a), x' \rangle$ where the limit $x' := \lim_{n \rightarrow \infty} \delta \circ (\pi_\lambda \odot \kappa_\eta)(b_n) \Omega_\lambda$ exists in H_λ by unitary equivalence between the cyclic representations $(H_\lambda, \delta \circ (\pi_\lambda \odot \kappa_\eta), \Omega_\lambda)$ and $(H_\odot, \pi_\omega(\cdot)|_{H_\odot}, \Omega_\omega)$ of $(F \odot \tilde{F}, \Delta_\lambda)$. Therefore $x' = 0$, and so $x = 0$ by the same unitary equivalence. Hence $H'_\lambda = H_\odot$. In a similar fashion $H_{\tilde{\lambda}} = H_\odot$ proving the claim. By unitary equivalence one can choose $(H_\lambda, \pi_\lambda, \Omega_\lambda)$ to be $(H'_\lambda, \pi_\mu(\cdot)|_{H'_\lambda}, \Omega_\mu)$, with a corresponding change in

κ_η , without changing the state Δ_λ . In conclusion we therefore have

$$H_{\tilde{\lambda}} = H_\lambda$$

when $\omega|_{F \odot \tilde{F}} = \Delta_\lambda$.

In this setting we now have the following result:

Proposition 3.6. *Suppose that ω is a coupling of (A, μ) and $(\tilde{B}, \tilde{\nu})$ such that $\omega|_{F \odot \tilde{F}} = \Delta_\lambda$. Then $\omega = \mu \odot_\lambda \tilde{\nu}$ if and only if any of the following three equivalent conditions are satisfied: $(H_\mu \ominus H_\lambda) \perp (H_{\tilde{\nu}} \ominus H_\lambda)$, $(H_\mu \ominus H_\lambda) \perp H_{\tilde{\nu}}$, or $H_\mu \perp (H_{\tilde{\nu}} \ominus H_\lambda)$.*

Proof. The equivalence of the three conditions is easily verified. So assume $\omega = \mu \odot_\lambda \tilde{\nu}$ and consider any $x \in H_\mu \ominus H_\lambda$ and $y \in H_{\tilde{\nu}}$, as well as sequences $a_n \in A$ and $b_n \in \tilde{B}$ such that $\gamma_\mu(a_n) \rightarrow x$ and $\gamma_{\tilde{\nu}}(b_n) \rightarrow y$. Then

$$\begin{aligned} \langle x, \gamma_{\tilde{\nu}}(b_n) \rangle &= \lim_{m \rightarrow \infty} \omega(a_m^* \otimes b_n) = \lim_{m \rightarrow \infty} \mu \odot_\lambda \tilde{\nu}(a_m^* \otimes b_n) \\ &= \left\langle \Omega_\lambda, (\pi_\lambda \circ D(a_m^*)) \left(\kappa_\eta \circ \tilde{D}(b_n) \right) \Omega_\lambda \right\rangle \\ &= \left\langle \Omega_\lambda, (\pi_\lambda \circ D(a_m^*)) \left(\kappa_\eta \circ \tilde{D} \left(\tilde{D}(b_n) \right) \right) \Omega_\lambda \right\rangle \\ &= \left\langle x, \gamma_{\tilde{\nu}} \left(\tilde{D}(b_n) \right) \right\rangle \\ &= 0 \end{aligned}$$

and therefore $\langle x, y \rangle = 0$. It follows that $(H_\mu \ominus H_\lambda) \perp H_{\tilde{\nu}}$. Conversely, assume that $(H_\mu \ominus H_\lambda) \perp (H_{\tilde{\nu}} \ominus H_\lambda)$ and consider any $a \in A$ and $b \in \tilde{B}$. Let P and \tilde{P} respectively be the projections of H_μ and $H_{\tilde{\nu}}$ on H_λ . Then for any $a \in A$ and $b \in \tilde{B}$

$$\begin{aligned} \omega(a \otimes b) &= \langle \gamma_\mu(a^*), \gamma_{\tilde{\nu}}(b) \rangle = \left\langle P\gamma_\mu(a^*), \tilde{P}\gamma_{\tilde{\nu}}(b) \right\rangle \\ &= \left\langle \gamma_\mu(D(a^*)), \gamma_{\tilde{\nu}}(\tilde{D}(b)) \right\rangle = \omega(D(a) \otimes \tilde{D}(b)) \\ &= \Delta_\lambda(D(a) \otimes \tilde{D}(b)) = \mu \odot_\lambda \tilde{\nu}(a \otimes b) \end{aligned}$$

which is sufficient. \square

We will also unitarily represent the dynamics of $\tilde{\mathbf{B}}$ on $H_{\tilde{\nu}}$, and we denote this representation by \tilde{V} . This representation is used in Sections 4 and 5 instead of V on H_ν to fit into the setting of Proposition 3.6.

4. RELATIVE ERGODICITY

In this section and the next we study how the relatively independent joinings constructed in the previous section relate to properties of systems. In particular in this section we consider relative ergodicity, which is a simple generalization of ergodicity. See for example [15, Section 6.6] for a discussion of the classical case.

The relevant terminology and definitions are as follows: The *fixed point algebra* of a system \mathbf{A} (or, put differently, of α) is defined as

$$A^\alpha := \{a \in A : \alpha_g(a) = a \text{ for all } g \in G\}$$

and this gives an identity subsystem \mathbf{A}^α of \mathbf{A} by simply restricting the state of \mathbf{A} to A^α . We call \mathbf{A}^α the *fixed point subsystem* of \mathbf{A} . Similarly, we call

$$H_\mu^U := \{x \in H_\mu : U_g x = x \text{ for all } g \in G\}$$

the fixed point space of U .

Definition 4.1. Let \mathbf{F} be a subsystem of \mathbf{A} . We call \mathbf{A} *ergodic relative to \mathbf{F}* if $A^\alpha \subset \zeta(F)$.

If \mathbf{F} in Definition 4.1 is the trivial system, then \mathbf{A} is simply called *ergodic*.

We note the following important facts:

Proposition 4.2. *The fixed point subsystem \mathbf{A}^α is a modular subsystem of \mathbf{A} . Furthermore, $H_\mu^U = \overline{\gamma_\mu(A^\alpha)}$.*

Proof. By [36, Corollary VIII.1.4] $\alpha_g \circ \sigma_t^\mu = \alpha_g \circ \sigma_t^{\mu \circ \alpha_g} = \sigma_t^\mu \circ \alpha_g$ from which $\sigma_t^\mu(A^\alpha) = A^\alpha$ follows.

We now show that $H_\mu^U = \overline{\gamma_\mu(A^\alpha)}$. It is clear that $\overline{\gamma_\mu(A^\alpha)} \subset H_\mu^U$. The converse follows from the Kovács-Szücs mean ergodic theorem [6, Proposition 4.3.8] (also see [21]) using an argument similar to that of [6, Theorem 4.3.20] (also see [18]): Let P be the projection of H_μ onto H_μ^U , then, since $\Omega_\mu \in H_\mu^U$ is cyclic for $\pi_\mu(A)'$, there exists a unique normal $(\pi_\mu \circ \alpha_g \circ \pi_\mu^{-1})_{g \in G}$ invariant projection $E : \pi_\mu(A) \rightarrow \pi_\mu(A^\alpha)$. Furthermore, E has the property that Ea is the unique element of $\pi_\mu(A)$ such that $(Ea)P = PaP$, for every $a \in \pi_\mu(A)$. Hence $(Ea)\Omega_\mu = Pa\Omega_\mu$ from which we obtain $\overline{\gamma_\mu(A^\alpha)} \supset H_\mu^U$. \square

Now for the main result of this section:

Theorem 4.3. *Let the identity system \mathbf{F} be a modular subsystem of \mathbf{A} . If \mathbf{A} and $\tilde{\mathbf{B}}$ are disjoint over \mathbf{F} for $\mathbf{B} = \mathbf{A}^\alpha$, then \mathbf{A} is ergodic relative to \mathbf{F} . On the other hand, if \mathbf{A} is ergodic relative to \mathbf{F} , then \mathbf{A} and $\tilde{\mathbf{B}}$ are disjoint over \mathbf{F} for every identity system \mathbf{B} which has \mathbf{F} as a modular subsystem.*

Proof. Without loss we assume that ζ is an inclusion, so $F \subset A$. Now, suppose \mathbf{A} is not ergodic relative to \mathbf{F} , in other words F is strictly contained in A^α . Note that \mathbf{F} is a modular subsystem of \mathbf{A}^α , since \mathbf{F} and \mathbf{A}^α are modular subsystems of \mathbf{A} . Then apply Proposition 3.5 with $\mathbf{B} = \mathbf{R} = \mathbf{A}^\alpha$ to obtain $\mu \odot_\psi \tilde{\nu} \neq \mu \odot_\lambda \tilde{\nu}$. But both these joinings are contained in $J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$ by Propositions 3.2 and 3.5, so \mathbf{A} and $\tilde{\mathbf{B}}$ are not disjoint over \mathbf{F} .

Conversely, assume that \mathbf{A} is ergodic relative to \mathbf{F} , and let \mathbf{B} be any identity system which has \mathbf{F} as a modular subsystem. We are going to apply Proposition 3.6. So consider any $\omega \in J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$. From this joining we obtain a conditional expectation operator $P_\omega : H_{\tilde{\nu}} \rightarrow H_\mu$ (i.e. $\langle x, P_\omega y \rangle = \langle x, y \rangle$ for all $x \in H_\mu$ and $y \in H_{\tilde{\nu}}$) such that $U_g P_\omega = P_\omega \tilde{V}_g$ (see [9, Proposition 2.4]). Then for any $b \in \tilde{B}$ we have $U_g P_\omega \gamma_{\tilde{\nu}}(b) = P_\omega \gamma_{\tilde{\nu}}(b)$, since $\tilde{\mathbf{B}}$ is an identity system. Since $F = A^\alpha$, it follows from Proposition 4.2 that $P_\omega \gamma_{\tilde{\nu}}(b) \in \overline{\gamma_\mu(A^\alpha)} = H_\lambda$. Noting that P_ω is the projection of H_ω onto H_μ restricted to $H_{\tilde{\nu}}$, we conclude that $H_{\tilde{\nu}} \perp (H_\mu \ominus H_\lambda)$, and hence $\omega = \mu \odot_\lambda \tilde{\nu}$ by Proposition 3.6. \square

This result contains [10, Theorem 2.1] as a special case by Proposition 3.4, namely \mathbf{A} is ergodic if and only if it is disjoint (over the trivial system) from all identity systems. Also note that in Theorem 4.3, \mathbf{A} being ergodic relative to \mathbf{F} means exactly that $A^\alpha = \zeta(F)$, i.e. \mathbf{F} is isomorphic to \mathbf{A}^α , since \mathbf{F} is assumed to be an identity system.

5. COMPACT SUBSYSTEMS

In this section we study an analogue of the previous section in terms of compact systems instead of identity systems. We therefore proceed by first building up some background regarding compact subsystems, which includes extending certain results from [26, Section 4] to more general group actions in a way (see Lemma 5.3 and Theorem 5.4 below) that is relevant for our intended applications. Our proofs of these extensions follow the basic plans of those in [26], but many of the details differ. (To avoid any confusion, we note that our terminology regarding topologies on $B(H)$ differs slightly from that of [26] and furthermore we do not always use the same topologies that [26] used in their results; we use the terminology of [6, Section 2.4.1], in particular the *weak* topology on $B(H)$ is generated by $\langle x, (\cdot)y \rangle$ for $x, y \in H$.) The additional definitions that we use in this section are the following:

Definition 5.1. A system \mathbf{A} is *compact* if the orbit $U_G x$ is totally bounded (i.e. $\overline{U_G x}$ is compact) in H_μ for every $x \in H_\mu$. An *eigenvector* of U is an $x \in H_\mu \setminus \{0\}$ such that there is a function, called its *eigenvalue*, $\chi : G \rightarrow \mathbb{C}$ such that $U_g x = \chi(g)x$ for all $g \in G$. The set of all eigenvalues is denoted by $\sigma_{\mathbf{A}}$. Let H_0 denote the Hilbert subspace of H_μ spanned by the eigenvectors of U . We will call $u \in A \setminus \{0\}$ an *eigenoperator* of α if there is a function $\chi : G \rightarrow \mathbb{C}$, its *eigenvalue*, such that $\alpha_g(u) = \chi(g)u$ for all $g \in G$.

As already mentioned, our immediate goal is to study compact subsystems of a given system. We start with an analogue of the first part of Proposition 4.2:

Proposition 5.2. For a system \mathbf{A} , denote by A^K the von Neumann subalgebra of A generated by the eigenoperators of α . Then $\alpha_g(A^K) =$

A^K , which allows us to define a subsystem $\mathbf{A}^K = (A^K, \mu^K, \alpha^K)$ of \mathbf{A} (in terms of an inclusion as the imbedding into \mathbf{A}) by setting $\mu^K := \mu|_{A^K}$ and $\alpha_g^K(a) := \alpha_g(a)$ for all $a \in A^K$. The system \mathbf{A}^K is a compact modular subsystem of \mathbf{A} .

Proof. Let S be the $*$ -algebra generated by the eigenoperators of α . Then clearly $\alpha_g(S) = S \subset A^K$. But $A^K = S''$, so S is σ -weakly dense in A^K , while α_g is σ -weakly continuous for each g . It follows that $\alpha_g(A^K) = A^K$.

Next we prove that the subsystem \mathbf{A}^K so obtained, is compact. Note that for any eigenoperator u of α , $\gamma_\mu(u)$ is an eigenvector of U . Furthermore, it is easily seen that if u and v are eigenoperators of α , then u^* is also eigenoperator of α , while uv is either zero or an eigenoperator of α . It follows that $\gamma_\mu(S) \subset H_0$, hence $\gamma_\mu(A^K) \subset H_0$. It is easy to show from the definition of H_0 that each of its elements has a totally bounded (i.e. relatively compact) orbit under U . Also note that (A^K, μ) can be cyclically represented on $H_\mu^K := \overline{\gamma_\mu(A^K)}$, which is contained in H_0 , with the unitary representation of α^K given by the restriction of U_g to H_μ^K . Therefore \mathbf{A}^K is indeed compact.

Lastly note that as in Proposition 4.2, for any eigenoperator u of α with eigenvalue χ , we have that $\theta_g(\sigma_t^\mu(u)) = \sigma_t^\mu(\theta_g(u)) = \chi(g)\sigma_t^\mu(u)$ so $\sigma_t^\mu(u) \in S$. It follows that $\sigma_t^\mu(S) = S$, hence $\sigma_t^\mu(A^K) = A^K$, i.e. \mathbf{A}^K is a modular subsystem of \mathbf{A} . \square

In order to go further, we first prove a technical result which is a version of [26, Lemma 4.1] appropriate for our needs:

Lemma 5.3. *Let H be a Hilbert space and T a weakly closed vector subspace of $B(H)$. Assume that there is a unit vector $\Omega \in H$ such that $\overline{T'\Omega} = H$. Let G be an arbitrary group and $W_g : H \rightarrow H$ a linear isometry for every $g \in G$ such that $W_g W_h = W_{gh}$ for all $g, h \in G$, and with*

$$K_a := \overline{\{W_g a \Omega : g \in G\}}$$

compact in H for every $a \in T$. Endow $B(T)$ with the topology of pointwise strong convergence generated by the seminorms $\Theta \mapsto \|\Theta(a)x\|$ where $a \in T$ and $x \in H$. Let Γ be the set of all linear contractions $\Theta \in B(T)$ such that $\Theta(a)\Omega \in K_a$ for all $a \in T$, and view it as a topological subspace of $B(T)$. Then Γ is a compact group with composition as its operation and the identity map on T as its identity.

Proof. Let us first consider the topology on Γ . For $a \in T, a' \in T', x \in H$ and $\Theta_1, \Theta_2 \in B(T)$ we have

$$\|\Theta_1(a)x - \Theta_2(a)x\| \leq 2\|a\| \|x - a'\Omega\| + \|a'\| \|\Theta_1(a)\Omega - \Theta_2(a)\Omega\|$$

from which (together with $\overline{T'\Omega} = H$) it follows that the topology on $B(T)$ is in fact generated by the seminorms $\Theta \mapsto \|\Theta(a)\Omega\|$ where $a \in T$,

hence we have a simplified description of the topology on Γ . Now consider the compact space

$$K := \prod_{a \in T} K_a$$

and the function

$$f : \Gamma \rightarrow K : \Theta \mapsto (\Theta(a)\Omega)_{a \in T}$$

which is easily shown to be injective because of $\overline{T'\Omega} = H$. From the simplified description of the topology on Γ described above and the definition of the product topology on K it is clear that $f : \Gamma \rightarrow f(\Gamma)$ is a homeomorphism. We now show that $f(\Gamma)$ is closed:

Let $x = (x(a))_{a \in T}$ be in the closure of $f(\Gamma)$. Let $(\Theta_i)_i$ be a net in Γ such that $(f(\Theta_i))_i$ converges to x , then from the inequality above $(\Theta_i(a)y)_i$ is a Cauchy net in H for every $y \in H$ and $a \in T$. This allows us to define $\Theta(a)y := \lim_i \Theta_i(a)y$ and since each Θ_i is a contraction, $\Theta(a) \in B(H)$ and $\|\Theta(a)\| \leq \|a\|$. By this definition $(\Theta_i(a))_i$ converges strongly and therefore also weakly to $\Theta(a)$, but T is weakly closed therefore $\Theta(a) \in T$. Since $\Theta_i(a)\Omega \in K_a$ and K_a is closed, it again follows from the definition of $\Theta(a)$ that $\Theta(a)\Omega \in K_a$ for all $a \in T$. This shows that $\Theta \in \Gamma$. But by our choice of $(\Theta_i)_i$ we have $x(a) = \lim_i \Theta_i(a)\Omega = \Theta(a)\Omega$ so $x = f(\Theta) \in f(\Gamma)$ and therefore $f(\Gamma)$ is indeed closed.

It follows that $f(\Gamma)$ is compact, since K is, and therefore Γ is compact. Also note that Γ is Hausdorff.

Now we start to prove the group structure on Γ . Note that $W_1 = 1$, since $W_g W_1 = W_g$ while W_g is injective. It follows that $\text{id}_T \in \Gamma$. For any $\Theta_1, \Theta_2 \in \Gamma$ we have

$$\Theta_1 \circ \Theta_2(a)\Omega \in K_{\Theta_2(a)} \subset \overline{\{W_g W_h a \Omega : g, h \in G\}} = K_a$$

and therefore $\Theta_1 \circ \Theta_2 \in \Gamma$.

Next we show the continuity of this product on Γ . For $\Theta \in \Gamma$ we have $\Theta(a)\Omega \in K_a \subset \{x \in H : \|x\| = \|a\Omega\|\}$, since each W_g is an isometry. Therefore we have a well-defined linear isometry

$$W_\Theta : \overline{T\Omega} \rightarrow \overline{T\Omega}$$

such that $W_\Theta(a\Omega) = \Theta(a)\Omega$, so in effect W_Θ represents Θ on $\overline{T\Omega}$. Using W_Θ one can show that

$$\begin{aligned} & \|(\Theta_1 \circ \Theta'_1)(a)\Omega - (\Theta_2 \circ \Theta'_2)(a)\Omega\| \\ & \leq \|\Theta'_1(a)\Omega - \Theta'_2(a)\Omega\| + \|\Theta_1(\Theta'_2(a))\Omega - \Theta_2(\Theta'_2(a))\Omega\| \end{aligned}$$

for all $\Theta_1, \Theta_2, \Theta'_1, \Theta'_2 \in \Gamma$ and $a \in T$. From this it follows that the product on Γ is indeed continuous.

We now complete the proof that Γ is a group by considering inverses. We have already shown that Γ is a semigroup, so for $\Theta \in \Gamma$

it follows that $\Theta^n \in \Gamma$ for every $n \in \mathbb{N} = \{1, 2, 3, \dots\}$. Furthermore $W_\Theta^n(a\Omega) = \Theta^n(a)\Omega \in K_a$, so since K_a is compact, the orbit $(W_\Theta^n(a\Omega))_{n \in \mathbb{N}}$ is relatively compact (and therefore totally bounded) for every $a \in T$. Since W_Θ is an isometry, it follows from [26, Corollary 9.10] (and remarks made just after it) that for every $\varepsilon > 0$ and every finite set $E \subset T$ there exists an $n(E, \varepsilon) \in \mathbb{N}$ such that

$$\|\Theta^{n(E, \varepsilon)}(a)\Omega - a\Omega\| < \varepsilon$$

for all $a \in E$. Defining $(E, \varepsilon) \leq (E', \varepsilon')$ to mean $E \subset E'$ and $\varepsilon \geq \varepsilon'$, we obtain a net $(\Theta^{n(E, \varepsilon)})_{(E, \varepsilon)}$ in Γ which is seen to converge to id_T by using the description of the topology given at the beginning of this proof. Since Γ is compact, the net $(\Theta^{n(E, \varepsilon)-1})_{(E, \varepsilon)}$ has a limit point in Γ , say Θ' , and it is then not too difficult to show that $\Theta \circ \Theta' = \text{id}_T = \Theta' \circ \Theta$. In other words every element of Γ has an inverse in Γ , and therefore Γ is a group.

By the single theorem appearing in [11], it follows that Γ is a topological group, which completes the proof. \square

Now we can state and prove the basic result regarding compact subsystems, which is the version of [26, Theorem 4.2] that we will use. In this result and the rest of the section our group G is assumed to be abelian, although this assumption is unnecessary in the first part of the proof of the next theorem, as will be indicated.

Theorem 5.4. *Consider a system \mathbf{A} with G assumed to be abelian. Set*

$$\begin{aligned} T &:= \{a \in \pi_\mu(A) : a\Omega_\mu \in H_0\} \\ \xi(a) &:= \langle \Omega_\mu, a\Omega_\mu \rangle \\ \tau_g(a) &:= U_g a U_g^* \end{aligned}$$

for all $a \in T$ and $g \in G$. Then $\mathbf{T} := (T, \xi, \tau)$ is a subsystem of \mathbf{A} which is isomorphic to \mathbf{A}^K with the isomorphism given by $\pi_\mu^{-1}|_T : T \rightarrow A^K$. It follows that \mathbf{A}^K is the largest compact subsystem of \mathbf{A} in the sense that A^K contains the image of the algebra of any compact subsystem \mathbf{F} of \mathbf{A} (under the imbedding of \mathbf{F} in \mathbf{A}). Furthermore, $\overline{T\Omega_\mu} = H_0$, so \mathbf{T} can be cyclically represented on H_0 and $\sigma_{\mathbf{T}} = \sigma_{\mathbf{A}}$.

Proof. The notation $M = \pi_\mu(A)$, $\theta_g : M \rightarrow M : a \mapsto U_g a U_g^*$, $H = H_\mu$ and $\Omega = \Omega_\mu$ will be used in what follows. We divide the proof into a number of parts.

(i) First we show (see also [26, Proposition 3.2]) that for any $\chi \in \sigma_{\mathbf{A}}$ we have

$$\overline{M_\chi \Omega} = H_\chi$$

where

$$M_\chi := \{a \in M : \theta_g(a) = \chi(g)a \text{ for all } g \in G\}$$

and

$$H_\chi := \{x \in H : U_g x = \chi(g)x \text{ for all } g \in G\}.$$

To do this, let P_χ be the projection of H on H_χ . Note that $|\chi(g)| = 1$ and $\chi(g)\chi(h) = \chi(gh)$. It follows that $g \mapsto \overline{\chi(g)}U_g$ is also a unitary group on H , and therefore by the Alaoglu-Birkhoff mean ergodic theorem [6, Proposition 4.3.4] (also see [1]) we know that P_χ is in the strong closure of the convex hull of $\{\overline{\chi(g)}U_g : g \in G\}$. Thus there is a net $(m_\iota)_\iota$ given by

$$m_\iota = \sum_{j=1}^{n_\iota} w_{\iota j} \overline{\chi(g_{\iota j})} U_{g_{\iota j}}$$

with $w_{\iota 1} + \dots + w_{\iota n_\iota} = 1$ and $w_{\iota j} \geq 0$, which converges strongly to P_χ .

Now for any $a \in M$ the net $(c_\iota)_\iota$ given by

$$c_\iota(a) = \sum_{j=1}^{n_\iota} w_{\iota j} \overline{\chi(g_{\iota j})} \theta_{g_{\iota j}}(a)$$

has a weakly convergent subnet, say $(c_{\iota(\kappa)}(a))_\kappa$, since the unit ball of $B(H)$ is weakly compact [6, Proposition 2.4.2]. Since M is a von Neumann algebra, it is weakly closed, so the limit of this subnet, say a_0 , is in M . Since $(m_{\iota(\kappa)})_\kappa$ converges strongly to P_χ and $U_g \Omega = \Omega$, we have that

$$\langle x, P_\chi a \Omega \rangle = \lim_\kappa \langle x, m_{\iota(\kappa)} a \Omega \rangle = \lim_\kappa \langle x, c_{\iota(\kappa)}(a) \Omega \rangle = \langle x, a_0 \Omega \rangle$$

for any $x \in H$, so $P_\chi a \Omega = a_0 \Omega$. It follows that $\theta_g(a_0) \Omega = \chi(g) a_0 \Omega$, but Ω is separating for M , hence $a_0 \in M_\chi$.

We remark that one can in fact go further: Taking $b \in M'$ we have

$$c_{\iota(\kappa)}(a) b \Omega = b m_{\iota(\kappa)} a \Omega \rightarrow b P_\chi a \Omega = a_0 b \Omega$$

since $c_{\iota(\kappa)}(a) \in M$. Since $M' \Omega$ is dense in H , it is now straightforward to show that $b \Omega$ can be replaced by any $x \in H$, in other words $(c_{\iota(\kappa)}(a))$ actually converges strongly to a_0 . We will only need the convergence $c_{\iota(\kappa)}(a) \Omega \rightarrow a_0 \Omega$, though.

Now consider any $x \in H_\chi$ and $\varepsilon > 0$, and let $a \in M$ be such that $\|a \Omega - x\| < \varepsilon$. Then

$$\left\| \overline{\chi(g)} \theta_g(a) \Omega - x \right\| = \|U_g a \Omega - \chi(g)x\| < \varepsilon$$

so $\|c_{\iota(\kappa)}(a) \Omega - x\| < \varepsilon$. But by the convergence above, there is a κ such that $\|a_0 \Omega - c_{\iota(\kappa)}(a) \Omega\| < \varepsilon$, so $\|a_0 \Omega - x\| < 2\varepsilon$. This proves that $M_\chi \Omega$ is dense in H_χ as required, since clearly $M_\chi \Omega \subset H_\chi$.

(ii) Now we prove a number of properties of T . First, since $U_g H_0 \subset H_0$ from the definition of H_0 while $U_g^* \Omega = \Omega$, we clearly have $\tau_g(T) = \theta_g(T) = T$. It follows that $\tau_g \in B(T)$.

Next, set

$$S := \text{span} \bigcup_{\chi \in \sigma_{\mathbf{A}}} M_{\chi}$$

where “span” means finite linear combinations. From (i) it follows that $\overline{S\Omega} = H_0$. It is also readily verified that S is a $*$ -algebra contained in T ; in fact, S is the $*$ -algebra generated by the eigenoperators of θ . In particular $\overline{T\Omega} = H_0$.

Lastly we show that T is weakly closed, so consider any a in the weak closure of T , and a net (a_{ι}) in T converging weakly to a . Since M is weakly closed, $a \in M$. Furthermore, from the definition of T , for every $x \in H_0^{\perp}$ we have $0 = \langle x, a_{\iota}\Omega \rangle \rightarrow \langle x, a\Omega \rangle$, so $a\Omega \in H_0$. Therefore, again from the definition of T , we have $a \in T$ as required.

(iii) Everything so far in the proof holds for non-abelian G as well, but in the rest of the proof we do make use of the fact that G is abelian, since we are going to work with its Bohr compactification.

In particular we now show that T is a von Neumann algebra using the group Γ given by Lemma 5.3 applied to T with $W = U$. Note that all the requirements in Lemma 5.3 are satisfied: $M' \subset T'$ so $T'\Omega$ is dense in H , and for every $a \in T$ the closure K_a of the corresponding orbit in H is indeed compact, since each element of H_0 has a totally bounded (i.e. relatively compact) orbit under U (in fact H_0 contains all such elements [5, Lemma 6.6], but this fact will only be used in (iv) below).

Assign the discrete topology to G . Since G is abelian, we can imbed it into its Bohr compactification \bar{G} . Let us denote this canonical imbedding by $\iota : G \rightarrow \bar{G}$. It is an injective group homomorphism with $\iota(G)$ dense in \bar{G} . The dual map of this is a group isomorphism $\hat{\iota} : \hat{\bar{G}} \rightarrow \hat{G}$. For $\chi \in \hat{\bar{G}}$, set $\langle \cdot, \chi \rangle := \hat{\iota}^{-1}(\chi)$, then it follows that $\langle \iota(g), \chi \rangle = \chi(g)$. Note that since $\tau_g \in B(T)$ is an isometry, it follows that $\tau : G \rightarrow \Gamma : g \mapsto \tau_g$ is a well-defined group homomorphism (and it is continuous, since G is discrete) which by the universal property of the Bohr compactification [13, Proposition (4.78)] can be extended to a continuous group homomorphism

$$\bar{\tau} : \bar{G} \rightarrow \Gamma : g \mapsto \bar{\tau}_g$$

which means that $\bar{\tau}_{\iota(g)} = \tau_g$ for all $g \in G$.

Since $\Gamma \rightarrow H : \Theta \mapsto \Theta(a)y$ is continuous, so is $\bar{G} \rightarrow H : g \mapsto \bar{\tau}_g(a)y$ for all $a \in T$ and $y \in H$. Denoting the normalized Haar measure on the compact group \bar{G} by m , we can for every $a \in T$ and $\chi \in \hat{\bar{G}}$ define a unique $a_{\chi} \in B(H)$ by requiring

$$\langle x, a_{\chi}y \rangle = \int_{\bar{G}} \overline{\langle g, \chi \rangle} \langle x, \bar{\tau}_g(a)y \rangle dm(g)$$

for all $x, y \in H$. If f is a weakly continuous linear functional on $B(H)$ it follows (see for example [8, Theorem V.3.9]) that it is a finite linear

combination of such $\langle x, (\cdot)y \rangle$ forms, hence

$$(1) \quad f(a_\chi) = \int_{\bar{G}} \overline{\langle g, \chi \rangle} f(\bar{\tau}_g(a)) dm(g)$$

so if f is zero on T we have $f(a_\chi) = 0$ and therefore by [8, Corollary V.3.12] $a_\chi \in T$ since T is weakly closed. For $g \in G$ it follows from the definition of a_χ that

$$\begin{aligned} \langle x, \tau_g(a_\chi)y \rangle &= \langle U_g^* x, a_\chi U_g^* y \rangle = \int_{\bar{G}} \overline{\langle h, \chi \rangle} \langle U_g^* x, \bar{\tau}_h(a) U_g^* y \rangle dm(h) \\ &= \int_{\bar{G}} \overline{\langle h, \chi \rangle} \langle x, \bar{\tau}_{\iota(g)h}(a)y \rangle dm(h) \\ &= \int_{\bar{G}} \overline{\langle \iota(g)^{-1}h, \chi \rangle} \langle x, \bar{\tau}_h(a)y \rangle dm(h) \\ &= \overline{\langle \iota(g)^{-1}, \chi \rangle} \int_{\bar{G}} \overline{\langle h, \chi \rangle} \langle x, \bar{\tau}_h(a)y \rangle dm(h) \\ &= \chi(g) \langle x, a_\chi y \rangle \end{aligned}$$

and therefore $\tau_g(a_\chi) = \chi(g)a_\chi$ which means that $a_\chi \in S$, since $a_\chi \in T \subset M$. We now use this result to show that S is weakly dense in T , so take any $a \in T$. Let f be any weakly continuous linear functional on $B(H)$ which is zero on the weak closure $\overline{a_{\bar{G}}}$ of the span of $\{a_\chi : \chi \in \hat{G}\}$. From (1) and the Plancherel theorem it follows that $g \mapsto f(\bar{\tau}_g(a))$ is zero in $L^2(\bar{G})$. But this function is continuous by arguments as above, so $f(\bar{\tau}_g(a)) = 0$ for all $g \in \bar{G}$, for if not, then f would be non-zero on an open neighbourhood of some $g \in \bar{G}$, contradicting the fact that the Haar measure of a non-empty open set is non-zero. In particular for $g = 1$ we find $f(a) = 0$. This means that $a \in \overline{a_{\bar{G}}}$ and therefore S is weakly dense in T . Since S is a $*$ -algebra containing the identity operator, we conclude that T is a von Neumann algebra by the von Neumann density theorem [6, Corollary 2.4.15].

(iv) By (ii) and (iii) we know that \mathbf{T} is indeed a subsystem of \mathbf{A} . It has $\sigma_{\mathbf{T}} = \sigma_{\mathbf{A}}$, since it can be cyclically represented on H_0 (because of $\overline{T\Omega_\mu} = H_0$) with the dynamics given by $U_g|_{H_0}$. Since S above is the $*$ -algebra generated by the eigenoperators of θ , we have $S \subset \pi_\mu(A^K) \subset T$. Hence $T = \pi_\mu(A^K)$, since $\pi_\mu(A^K)$ is a von Neumann algebra and therefore weakly closed. From this the isomorphism between \mathbf{T} and \mathbf{A}^K follows. If \mathbf{F} is any compact subsystem of \mathbf{A} , and $a \in F$, then by definition of compactness of a system the orbit of $\gamma_\mu(\zeta(a))$ under U is totally bounded (keep in mind that (F, λ) can be cyclically represented on $\overline{\gamma_\mu(\zeta(F))}$, with the corresponding unitary representation of φ given by the restriction of U_g to this space), thus $\pi_\mu(\zeta(a))\Omega_\mu \in H_0$ according to [5, Lemma 6.6]. In other words $\pi_\mu(\zeta(a)) \in T$, and therefore $\zeta(a) \in A^K$ which means that \mathbf{A}^K is the largest compact subsystem of \mathbf{A} . \square

Note that the sole reason we needed to assume that G is abelian in this theorem, is that we use its Bohr compactification in the proof.

Whereas \mathbf{A}^α is by definition the largest identity subsystem of \mathbf{A} , we have seen above that \mathbf{A}^K is the largest compact subsystem of \mathbf{A} , at least when G is abelian. Clearly $A^\alpha \subset A^K$. We mention that a very simple version of this in the context of noncommutative topological dynamics was also discussed in [23, Section 2] and [4, Definition 1.2]. In analogy to Proposition 4.2, Theorem 5.4 says that $H_0 = \overline{\gamma_\mu(A^K)}$.

Before we return to relatively independent joinings we note the following generalization of [5, Theorem 6.8] (also see [26, Proposition 5.4]). Note that in terms of Definition 5.1's notation, we call \mathbf{A} *weakly mixing* if $\dim H_0 = 1$.

Corollary 5.5. *A system \mathbf{A} with G assumed abelian, is weakly mixing if and only if \mathbf{A}^K is the trivial system.*

Proof. If \mathbf{A}^K is trivial, then so is \mathbf{T} in Theorem 5.4, so $H_0 = \overline{T\Omega_\mu} = \mathbb{C}\Omega_\mu$. Conversely, if \mathbf{A} is weakly mixing, then $\overline{T\Omega_\mu} = \mathbb{C}\Omega_\mu$, but Ω_μ is separating for $\pi_\mu(A)$ and therefore for T , so $T = \mathbb{C}$, hence \mathbf{A}^K is trivial. \square

In addition to the assumptions in this corollary, in [5, Theorem 6.8] it was also assumed that \mathbf{A} is ergodic.

Now we finally return to relatively independent joinings, namely an analogue of Theorem 4.3. The proof is very similar to that of Theorem 4.3.

Theorem 5.6. *Assume that G is abelian, and let \mathbf{F} be a compact modular subsystem of \mathbf{A} . If \mathbf{A} and $\tilde{\mathbf{B}}$ are disjoint over \mathbf{F} for $\mathbf{B} = \mathbf{A}^K$, then \mathbf{F} is isomorphic to \mathbf{A}^K . On the other hand, if \mathbf{F} is isomorphic to \mathbf{A}^K , then \mathbf{A} and $\tilde{\mathbf{B}}$ are disjoint over \mathbf{F} for every compact system \mathbf{B} which has \mathbf{F} as a modular subsystem.*

Proof. Assume without loss that \mathbf{F} is imbedded in \mathbf{A} by inclusion. Suppose \mathbf{F} is not isomorphic to \mathbf{A}^K , then \mathbf{F} is strictly contained in \mathbf{A}^K by Theorem 5.4. Note that \mathbf{F} is a modular subsystem of \mathbf{A}^K , since \mathbf{F} and \mathbf{A}^K are modular subsystems of \mathbf{A} . Now apply Proposition 3.5 with $\mathbf{B} = \mathbf{R} = \mathbf{A}^K$ to obtain $\mu \odot_\psi \tilde{\nu} \neq \mu \odot_\lambda \tilde{\nu}$. But both these joinings are contained in $J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$ by Propositions 3.2 and 3.5, so \mathbf{A} and $\tilde{\mathbf{B}}$ are not disjoint over \mathbf{F} .

Conversely, assume that \mathbf{F} is isomorphic (and therefore equal) to \mathbf{A}^K . Consider any $\omega \in J_\lambda(\mathbf{A}, \tilde{\mathbf{B}})$. For any eigenvector $y \in H_{\tilde{\nu}}$ of \tilde{V} with eigenvalue $\chi \in \sigma_{\tilde{\mathbf{B}}}$ we have

$$U_g P_\omega y = P_\omega \tilde{V}_g y = \chi(g) P_\omega y$$

with P_ω as in the proof of Theorem 4.3. So $P_\omega y = 0$ or $\chi \in \sigma_{\mathbf{A}}$. Therefore $P_\omega y \in H_0 = H_\lambda$ with H_0 as in Definition 5.1 and H_λ as given

in the setting described before Proposition 3.6. Since $\tilde{\mathbf{B}}$ has discrete spectrum (i.e. $H_{\tilde{\nu}}$ is spanned by the eigenvectors of \tilde{V}) as remarked in [10, Proposition 2.6] (also see [22, Section 2.4] for the general theory), it follows that $P_{\omega}y \in H_{\lambda}$ for all $y \in H_{\tilde{\nu}}$. Since P_{ω} is the projection of H_{ω} onto H_{μ} restricted to $H_{\tilde{\nu}}$, it follows that $H_{\tilde{\nu}} \perp (H_{\mu} \ominus H_{\lambda})$, and hence $\omega = \mu \odot_{\lambda} \tilde{\nu}$ by Proposition 3.6. \square

In the case of trivial \mathbf{F} this theorem gives a variation on [10, Theorems 2.7 and 2.8] in that the latter assumed ergodicity of the systems involved, but did not require G to be abelian, namely:

Corollary 5.7. *Assuming G is abelian, \mathbf{A} is weakly mixing if and only if it is disjoint from all compact systems.*

Proof. Use trivial \mathbf{F} in Theorem 5.6, together with Proposition 3.4 and Corollary 5.5. \square

As in Section 4, these results illustrate how relatively independent joinings relate to the structure of W^* -dynamical systems, in particular to certain types of subsystems.

6. ERGODICITY AND COMPACTNESS

In this concluding section we turn away from joinings and focus on subsystems. The goal is to study \mathbf{A}^K , which was defined in Proposition 5.2, when \mathbf{A} is ergodic, i.e. when \mathbf{A}^{α} is trivial. This is related to, and largely motivated by [19, Theorem 1.3] (see Corollary 6.2 below).

Remember that a state μ on a von Neumann algebra A is called *tracial* if $\mu(ab) = \mu(ba)$ for all $a, b \in A$. The following result strengthens the fact that \mathbf{A}^K is a modular subsystem of \mathbf{A} when the latter is ergodic, and note that in this result we need not assume G is abelian:

Theorem 6.1. *Let \mathbf{A} be ergodic. Then A^K is contained in the fixed point algebra of the modular group σ^{μ} of μ . Furthermore, μ^K is then tracial, and A^K therefore a finite von Neumann algebra.*

Proof. Let u be any eigenoperator of α , with eigenvalue χ . Then as in the proof of Proposition 4.2

$$\alpha_g(\sigma_t^{\mu}(u)) = \sigma_t^{\mu}(\alpha_g(u)) = \chi(g)\sigma_t^{\mu}(u).$$

Since \mathbf{A} is ergodic and u and $\sigma_t^{\mu}(u)$ are eigenoperators with the same norm and the same eigenvalue, it follows from [34, Lemma 2.1(3)] that there is a number, say $\Lambda(t) \in \mathbb{C}$, with $|\Lambda(t)| = 1$, such that

$$\sigma_t^{\mu}(u) = \Lambda(t)u$$

for all $t \in \mathbb{R}$. Now we show that $\sigma_t^{\mu}(u) = u$ (also see the proof of [10, Lemma 5.2]). It follows from the group property of σ^{μ} that $\Lambda(s+t) = \Lambda(s)\Lambda(t)$ for all $s, t \in \mathbb{R}$, and since $t \mapsto \langle x, \pi_{\mu}(\sigma_t^{\mu}(u))y \rangle$ is continuous for all $x, y \in H_{\mu}$, it follows that $t \mapsto \Lambda(t)$ is continuous. Therefore

$$\Lambda(t) = e^{i\theta t}$$

for all $t \in \mathbb{R}$ for some $\theta \in \mathbb{R}$; see for example [29, p. 12]. Denoting the modular operator associated with $(\pi_\mu(A), \Omega_\mu)$ by Δ , it follows that $\Delta^{it}\pi_\mu(u)\Omega_\mu = \pi_\mu(\sigma_t^\mu(u))\Omega_\mu = e^{i\theta t}\pi_\mu(u)\Omega_\mu$, hence by the definition of $J_\mu\Delta^{1/2}$ (see for example [6, Section 2.5.2])

$$J_\mu\pi_\mu(u)^*\Omega_\mu = J_\mu(J_\mu\Delta^{1/2})\pi_\mu(u)\Omega_\mu = \Delta^{1/2}\pi_\mu(u)\Omega_\mu = e^{\theta/2}\pi_\mu(u)\Omega_\mu$$

and by taking the norm both sides we conclude that $e^{\theta/2} = 1$, since $\alpha_g(uu^*) = |\chi(g)|^2 uu^* = uu^*$ and $\alpha_g(u^*u) = u^*u$, implying $uu^* = u^*u \in \mathbb{C}1 \setminus \{0\}$ by ergodicity and the fact that uu^* and u^*u are both positive. Therefore $\theta = 0$. This proves that

$$\sigma_t^\mu(u) = u$$

for all $t \in \mathbb{R}$. Then by the definition of A^K in Proposition 5.2, it is indeed contained in the fixed point algebra of the modular group σ^μ , since the latter algebra is a von Neumann algebra containing all eigenoperators of α , as just shown.

Since \mathbf{A}^K is a modular subsystem of \mathbf{A} , the modular group σ^{μ^K} of μ^K is simply given by the restriction of σ_t^μ to A^K . Hence $\sigma_t^{\mu^K} = \text{id}_{A^K}$ for all t by what we have shown above. It follows that the modular operator associated with the cyclic representation of (A^K, μ^K) is the identity operator. Therefore by the basic definitions of Tomita-Takesaki theory, and writing the cyclic representation of (A^K, μ^K) as (H, π, Ω) and the corresponding modular conjugation as J , we have $Ja\Omega = a^*\Omega$ for all $a \in \pi(A^K)$. From this it follows that μ^K is tracial, namely for any $a, b \in A^K$ we have

$$\begin{aligned} \mu(ab) &= \langle \Omega, \pi(ab)\Omega \rangle = \langle \pi(a)^*\Omega, \pi(b)\Omega \rangle = \langle J\pi(a)\Omega, \pi(b)\Omega \rangle \\ &= \langle J\pi(b)\Omega, \pi(a)\Omega \rangle = \langle \Omega, \pi(ba)\Omega \rangle \\ &= \mu(ba) \end{aligned}$$

and the existence of a faithful (numerical) trace on A^K implies that it is a finite von Neumann algebra (see for example [20, p. 505]). \square

To clarify the connection with [19, Theorem 1.3], we deduce it as a corollary from this theorem (but we only state it for an inverse temperature of -1):

Corollary 6.2. *Suppose that the dynamics α of \mathbf{A} is given by the modular group σ^μ , i.e. $G = \mathbb{R}$ and $\alpha_t = \sigma_t^\mu$ for all $t \in \mathbb{R}$. If \mathbf{A} is ergodic, it then follows that it is weakly mixing.*

Proof. By Theorem 6.1 $A^K \subset A^\alpha$, but $A^\alpha = \mathbb{C}$, since \mathbf{A} is ergodic, therefore $A^K = \mathbb{C}$. Since \mathbb{R} is an abelian group, it follows from Corollary 5.5 that \mathbf{A} is weakly mixing. \square

One may wonder whether more generally, when \mathbf{A}^α is not necessarily trivial, one has $A^K = A^\alpha$ if α is the modular group of μ . However on

$A = B(H)$ with H a finite dimensional Hilbert space, one has the simple counter example given by the Gibbs state (at inverse temperature -1) $\mu(a) = \text{Tr}(e^h a) / \text{Tr}(e^h)$ where $h \in A$ is hermitian, in which case $\alpha_t(a) = e^{iht} a e^{-iht}$ is the modular group of μ ; the system \mathbf{A} so obtained is compact but not in general an identity system, so $A^K = A \neq A^\alpha$.

Another corollary of Theorem 6.1 is the following (also see [27, Proposition 7.2]):

Corollary 6.3. *If a system \mathbf{A} is ergodic and compact, and G is abelian, then μ is necessarily tracial and A a finite von Neumann algebra.*

Proof. By Theorem 5.4 $\mathbf{A}^K = \mathbf{A}$, and the result then follows from Theorem 6.1. \square

Results of the nature of Corollary 6.3 have of course been well studied. However, our assumption that the system is compact makes Corollary 6.3 a relatively easy result. In the literature much more difficult results have been obtained; see [17, 32, 33, 34, 24, 2, 27, 16] for the development.

The results of this section illustrate how subsystems can be applied to derive interesting properties of systems, and show how aspects of the work above fit into previous literature on noncommutative dynamical systems.

ACKNOWLEDGMENTS

I thank Anton Ströh for useful discussions, and Richard de Beer for pointing the paper [19] out to me long ago. This work was supported by the National Research Foundation of South Africa.

REFERENCES

- [1] L. Alaoglu and G. Birkhoff, *General ergodic theorems*, Ann. of Math. (2) 41, (1940). 293–309. MR0002026 (1,339a)
- [2] S. Albeverio and R. Høegh-Krohn, *Ergodic actions by compact groups on C^* -algebras*, Math. Z. **174** (1980), 1–17. MR0591609 (82h:46087)
- [3] T. Austin, T. Eisner and T. Tao, *Nonconventional ergodic averages and multiple recurrence for von Neumann dynamical systems*, arXiv:0912.5093v4 [math.OA].
- [4] D. Avitzour, *Noncommutative topological dynamics. II*, Trans. Amer. Math. Soc. **282** (1984), 121–135. MR 0728705 (85i:46089)
- [5] C. Beyers, R. Duvenhage and A. Ströh, *The Szemerédi property in ergodic W^* -dynamical systems*, J. Operator Theory, to appear, arXiv:0709.1557v1 [math.OA].
- [6] O. Bratteli and D. W. Robinson, *Operator algebras and quantum statistical mechanics 1*, second edition, Springer-Verlag, New York, 1987. MR 0887100 (88d:46105)
- [7] O. Bratteli and D. W. Robinson, *Operator algebras and quantum statistical mechanics 2*, second edition, Springer-Verlag, Berlin, 1997. MR1441540 (98e:82004)

- [8] N. Dunford and J. T. Schwartz, *Linear operators. Part I. General theory*, With the assistance of William G. Bade and Robert G. Bartle. Reprint of the 1958 original. Wiley Classics Library. A Wiley-Interscience Publication. John Wiley & Sons, Inc., New York, 1988. MR1009162 (90g:47001a)
- [9] R. Duvenhage, *Joinings of W^* -dynamical systems*, J. Math. Anal. Appl. **343** (2008), 175–181. arXiv:0802.0827v1 [math.OA]. MR 2409467 (2009b:46135)
- [10] R. Duvenhage, *Ergodicity and mixing of W^* -dynamical systems in terms of joinings*, arXiv:0803.2147 [math.OA].
- [11] R. Ellis, *A note on the continuity of the inverse*, Proc. Amer. Math. Soc. **8** (1957), 372–373. MR0083681 (18,745d)
- [12] F. Fidaleo, *An ergodic theorem for quantum diagonal measures*, Infin. Dimens. Anal. Quantum Probab. Relat. Top. **12** (2009), 307–320. MR 2541399
- [13] G. B. Folland, *A course in abstract harmonic analysis*, Studies in Advanced Mathematics. CRC Press, Boca Raton, FL, 1995.
- [14] H. Furstenberg, *Disjointness in ergodic theory, minimal sets, and a problem in Diophantine approximation*, Math. Systems Theory **1** (1967), 1–49. MR 0213508 (35 #4369)
- [15] E. Glasner, *Ergodic theory via joinings*, Mathematical Surveys and Monographs 101, American Mathematical Society, Providence, RI, 2003. MR 1958753 (2004c:37011)
- [16] R. Høegh-Krohn, M. B. Landstad and E. Størmer, *Compact ergodic groups of automorphisms*, Ann. of Math. (2) **114** (1981), 75–86. MR0625345 (82i:46097)
- [17] N. M. Hugenholtz, *On the factor type of equilibrium states in quantum statistical mechanics*, Comm. Math. Phys. **6** (1967), 189–193. MR0225177 (37 #772)
- [18] A. Z. Jadczyk, *On some groups of automorphisms of von Neumann algebras with cyclic and separating vector*, Comm. Math. Phys. **13** (1969) 142–153. MR0264936 (41 #9525)
- [19] V. Jakšić and C.-A. Pillet, *A note on eigenvalues of Liouvilleans*, J. Statist. Phys. **105** (2001), 937–941. MR1869571 (2002j:46085)
- [20] R. V. Kadison and J. R. Ringrose, John R. *Fundamentals of the theory of operator algebras. Vol. II. Advanced theory*, Corrected reprint of the 1986 original. Graduate Studies in Mathematics, 16. American Mathematical Society, Providence, RI, 1997.
- [21] I. Kovács and J. Szücs, *Ergodic type theorems in von Neumann algebras*, Acta Sci. Math. (Szeged) **27** (1966) 233–246. MR0209857 (35 #753)
- [22] U. Krengel, *Ergodic theorems*, Walter de Gruyter & Co., Berlin, 1985. MR 0797411 (87i:28001)
- [23] D. Laisson and G. Laisson, *Topological dynamics on C^* -algebras*, Trans. Amer. Math. Soc. **204** (1975), 197–205. MR0358365 (50 #10831)
- [24] R. Longo, *Notes on algebraic invariants for noncommutative dynamical systems*, Comm. Math. Phys. **69** (1979), 195–207. MR0550019 (80j:46108)
- [25] S. Neshveyev and E. Størmer, *Dynamical entropy in operator algebras*, Ergebnisse der Mathematik und ihrer Grenzgebiete, 3. Folge, A Series of Modern Surveys in Mathematics, Vol. 50, Springer-Verlag, Berlin, 2006. MR2251116 (2007m:46108)
- [26] C. P. Niculescu, A. Ströh and L. Zsidó, *Noncommutative extensions of classical and multiple recurrence theorems*, J. Operator Theory **50** (2003), 3–52. MR 2015017 (2004k:46123)
- [27] D. Olesen, G. K. Pedersen and M. Takesaki, *Ergodic actions of compact abelian groups*, J. Operator Theory **3** (1980), 237–269.

- [28] S. Popa, *Cocycle and orbit equivalence superrigidity for malleable actions of w -rigid groups*, Invent. Math. **170** (2007), no. 2, 243–295.
- [29] W. Rudin, *Fourier analysis on groups*, Interscience Tracts in Pure and Applied Mathematics, No. 12 Interscience Publishers (a division of John Wiley and Sons), New York-London, 1962. MR0152834 (27 #2808)
- [30] D. J. Rudolph, *An example of a measure preserving map with minimal self-joinings, and applications*, J. Analyse Math. **35** (1979), 97–122. MR 0555301 (81e:28011)
- [31] J.-L. Sauvageot and J.-P. Thouvenot, *Une nouvelle définition de l'entropie dynamique des systèmes non commutatifs*, Comm. Math. Phys. **145** (1992), 411–423. MR 1162806 (93d:46113)
- [32] E. Størmer, *Types of von Neumann algebras associated with extremal invariant states*, Comm. Math. Phys. **6** (1967), 192–204. MR0225178 (37 #773)
- [33] E. Størmer, *Automorphisms and invariant states of operator algebras*, Acta Math. **127** (1971), 1–9. MR0383097 (52 #3978)
- [34] E. Størmer, *Spectra of ergodic transformations*, J. Funct. Anal. **15** (1974), 202–215. MR 0377544 (51 #13715)
- [35] Ş. Strătilă, *Modular theory in operator algebras*, Editura Academiei Republicii Socialiste România, Bucharest, Abacus Press, Tunbridge Wells, 1981. MR 0696172 (85g:46072)
- [36] M. Takesaki, *Theory of operator algebras. II*, Encyclopaedia of Mathematical Sciences, 125. Operator Algebras and Non-commutative Geometry, 6. Springer-Verlag, Berlin, 2003. MR1943006 (2004g:46079)

DEPARTMENT OF PHYSICS, UNIVERSITY OF PRETORIA, PRETORIA 0002, SOUTH AFRICA

E-mail address: rocco.duvenhage@up.ac.za